

Pulse-like and crack-like ruptures in experiments mimicking crustal earthquakes

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Theoretical studies have shown that the issue of rupture modes has important implications for fault constitutive laws, stress conditions on faults, energy partition and heat generation during earthquakes, scaling laws, and spatiotemporal complexity of fault slip. Early theoretical models treated earthquakes as crack-like ruptures, but seismic inversions indicate that earthquake ruptures may propagate in a self-healing pulse-like mode. A number of explanations for the existence of slip pulses have been proposed and continue to be vigorously debated. This study presents experimental observations of spontaneous pulse-like ruptures in a homogeneous linear-elastic setting that mimics crustal earthquakes; reveals how different rupture modes are selected based on the level of fault prestress; demonstrates that both rupture modes can transition to supershear speeds; and advocates, based on comparison with theoretical studies, the importance of velocity-weakening friction for earthquake dynamics.

earthquake physics | mechanics of faulting | rupture modes | shear cracks | velocity-dependent friction

Destructive large earthquakes occur as dynamic frictional ruptures along preexisting interfaces (or faults) in the Earth's crust. Inversions of seismic and other field observations have significantly advanced our understanding of earthquake ruptures. At the same time, detailed inversions often are impossible due to limited data availability or limited knowledge of the structure and properties of the crust. Numerical modeling of earthquakes helps pinpoint potential rupture scenarios but, in turn, requires a number of still poorly known inputs that are being actively researched. Such inputs include fault friction laws and initial stress conditions.

This reality highlights the need for highly instrumented laboratory experiments that reproduce some of the basic physics governing rupture dynamics of crustal earthquakes while preserving enough simplicity so that clear conclusions can be obtained by direct observation. One example of such experiments is the work of Xia *et al.* (1, 2), which has demonstrated the ability of the spontaneous dynamic ruptures to transition from sub-Rayleigh to supershear speeds. In the present study, we use their experimental configuration to investigate conditions leading to the selection of pulse-like vs. crack-like rupture modes in a setting that mimics crustal earthquakes. The geometry, loading, and nucleation mechanism are essentially 2D. This configuration is relevant for understanding the dynamics of large strike-slip earthquakes, which are dominated by in-plane sliding, and constitutes an experimental equivalent of 2D in-plane or Mode II numerical models of dynamic rupture, which are common in earthquake studies (3–12).

Numerical simulations in models that involve homogeneous elastic and interface properties and velocity-independent fault strength result in the crack-like mode of earthquake propagation, in which the duration of slip at each point on the fault is comparable to the overall rupture duration (3, 13–16). However, seismic inversions indicate (17) that ruptures on real faults may propagate in the pulse-like mode, in which slip duration at a point is much shorter than the overall rupture duration. Theo-

retical and numerical studies have shown that the issue of rupture modes may have important implications for fault constitutive laws, stress conditions on faults, energy partition and heat generation during earthquakes, scaling laws, and the spatiotemporal complexity of slip (5, 8, 12, 17–24). Pulse-like ruptures have been obtained in a number of numerical simulations that include significant weakening of interface friction with sliding velocity (5, 12, 18–22, 24). The simulations imply that fault friction may be characterized by significant velocity weakening, a conclusion further supported by a number of recent rock experiments and theoretical studies that have uncovered strongly velocity-weakening friction at seismic slip velocities (25–28). Other explanations for the occurrence of pulse-like ruptures include interaction of rupture with fault geometry or local heterogeneities (29–33) and normal stress variation due to differences in material properties across the interface (bimaterial effect) (6–8). Which mechanism dominates in real earthquakes remains an open research question.

In our experiments, there are no heterogeneities in interface properties or prestress and no bimaterial effect. Our goal is to determine whether pulse-like ruptures can occur in such a homogeneous configuration and, if so, what controls the selection of rupture modes. The only prior experimental study of different rupture modes under similarly homogeneous conditions was done with strong impact loading and interfaces with no shear prestress (34). However, those loading conditions are quite different from the ones on tectonically loaded faults in the Earth's crust. In the present study, we use an experimental configuration with an interface prestressed both in compression and in shear (1), simulating a tectonically loaded fault, and combine it with experimental diagnostics that let us conclusively determine the mode and speed of rupture propagation (34, 35). By systematically varying loading parameters, we observe pulse-like and crack-like ruptures, and systematic transition between them. Our results are consistent with the theoretical study of Zheng and Rice (21) who considered velocity-weakening interfaces and showed that selection of rupture modes depends on fault prestress and velocity-weakening properties of faults. We also present experimental observations of pulse-like ruptures transitioning to supershear speeds.

Experimental Design

The Earth's crust was simulated by a square (150 × 150 mm) photoelastic plate with a thickness of 9.5 mm. The plate was cut into two identical quadrilaterals, introducing an interface inclined at an angle α to the horizontal (Fig. 1). Care was taken to polish and then roughen the surfaces of the cut for each sample

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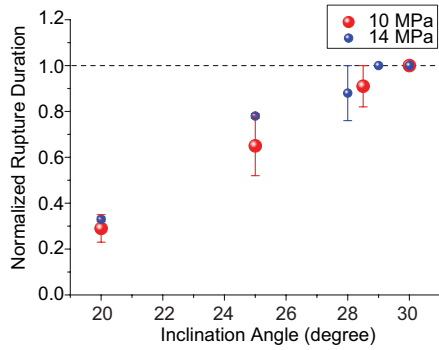


Fig. 3. Rupture duration, normalized by the maximum potential sliding time, plotted as a function of inclination angle α for two values of far-field pressure P . Normalized rupture duration of 1 corresponds to crack-like ruptures, whereas smaller values correspond to progressively narrower pulse-like ruptures.

exhibits a wider pulse compared with the previous case. Eq. 1 gives $18.2 \mu\text{s}$ as the time of sliding initiation, a time that is consistent with the arrival of the photoelastic fringe concentration. The two criteria of interface locking result in locking times of 42.6 and $54.6 \mu\text{s}$, respectively. These times correspond to a rupture duration between 24.4 and $36.4 \mu\text{s}$. Fig. 2d shows the superposition of the relative velocity record and photoelastic image for this case and demonstrates that the pulse is closer to a crack-like mode in the sense that locking occurs near the nucleation site.

For $\alpha = 30^\circ$, the prestress $f_0 = 0.577$ is close to the static coefficient of friction, $f_s = 0.6$, for the interface. In this case, the interface is close to the state of uniform sliding along the entire fault. Measurements presented in Fig. 2e and f reveal that the sliding mode is now fully crack-like in the sense that sliding at the measurement location does not stop during the observation time. The average relative velocity is ≈ 2 m/s (the maximum relative velocity is 7.35 m/s). A small “precursor” pulse precedes the arrival of both the shear wave front and the crack-like rupture that follows it and shows that the rupture begins to transition to supershear speeds close to the location of the relative velocity measurement. The details of such transitions are discussed in *Transition of Pulse-Like and Crack-Like Modes to Supershear Speeds*.

Collective Analysis of Rupture Modes

The three cases described in the preceding section were selected to demonstrate the existence of both pulse-like and crack-like ruptures in the experiments and to showcase the systematic transition of sub-Rayleigh rupture modes from pulse-like to crack-like under increasing shear prestress. In addition to these three cases, a number of other experiments were conducted to examine the phenomenon in a more complete angle and prestress spectrum and to test the reliability of the results. Fig. 3 displays rupture durations for a collection of these experimental observations. In Fig. 3, rupture duration for each experiment is normalized by the maximum potential sliding time, which is equal to the time window of observation minus the arrival time of the rupture front. This normalized rupture duration ranges from 0 (no sliding at the measurement location) to 1 (continued sliding from the rupture arrival to the end of the observation). Crack-like ruptures correspond to the normalized rupture duration of 1, with smaller values indicating pulse-like ruptures of progressively shorter duration. For each experiment, the ends of the interval correspond to two estimates of the rupture duration, with a filled dot giving the average value.

Fig. 3 clearly shows that the inclination angle α and, hence, the prestress f_0 are the dominant factors in determining the rupture mode, with smaller angles and, hence, lower prestress favoring pulse-like ruptures of shorter duration. For angles $\alpha = 20^\circ$ and $\alpha = 25^\circ$, the rupture is clearly pulse-like. Above $\alpha = 28^\circ$, the more conservative estimate of rupture duration gives the normalized rupture duration of 1, indicating that those ruptures may be fully crack-like. For angles $\alpha = 29^\circ$ and $\alpha = 30^\circ$, the ruptures are clearly crack-like. Fig. 3 presents results for inclination angles between 20° and 30° . Angles below 20° were not studied, but those experiments would have likely produced dying pulses or no sliding at the measurement location. For angles larger than 31° , prestress $f_0 = \tau_0/\sigma_0 = \tan\alpha$ would exceed the static friction coefficient of 0.6 , which would cause the sliding to occur over the entire interface at once.

The systematic transition of rupture modes from pulse-like to crack-like presented in this work is qualitatively consistent with the theoretical study of velocity-weakening interfaces by Zheng and Rice (21). Their analysis emphasized the determining role of velocity-weakening friction $\tau_{ss}(V)$ in promoting either pulse-like or crack-like behavior depending on the level of prestress τ_0 . If τ_{pulse} is the maximum value of τ_0 that satisfies $\tau_{\text{el}} = \tau_0 - \mu V/(2c_s) \leq \tau_{ss}(V)$ for all $V \geq 0$, then no crack-like solutions exist for $\tau_0 < \tau_{\text{pulse}}$. (Note that τ_{el} gives the elastodynamic stress for the case of uniform sliding along the entire interface.) For larger values of τ_0 , parameter T is defined as follows. The values of slip velocity V are found that solve $\tau_{\text{el}}(V) = \tau_{ss}(V)$. Let us denote the larger of the two possible solutions by V_{dyna} . Then

$$T = \frac{d\tau_{ss}/dV}{d\tau_{\text{el}}/dV} \bigg|_{V=V_{\text{dyna}}} \quad [2]$$

That is, parameter T is the ratio of the slopes of the steady-state friction curve τ_{ss} and elastodynamic stress τ_{el} evaluated at their intersection $V = V_{\text{dyna}}$. When T exists, it is a nondimensional scalar between zero and one. If T is close to zero, the rupture mode is crack-like. If T is close to one, the rupture mode is pulse-like. If T does not exist, the rupture mode is either pulse-like or there is no rupture propagation.

To apply the analysis of Zheng and Rice (21) to our experiments, we describe the frictional properties of the interface by Dieterich–Ruina rate-and-state friction law (38–43) enhanced with additional velocity weakening at high slip velocities, as appropriate for flash heating (28). For steady-state sliding, the friction law reduces to

$$\tau_{ss}(V) = \sigma \left(f_w + \frac{f_* + (a - b) \ln(V/V_*) - f_w}{1 + V/V_w} \right), \quad [3]$$

where f_* and V_* are, respectively, the reference friction coefficient and slip velocity, a and b are rate-and-state friction coefficients, V_w is the characteristic slip velocity for flash heating, and f_w is the residual friction coefficient at high sliding rates. Based on previous studies of frictional phenomena on Homalite interfaces (1, 44), we use the following values: $f_* = 0.6$, $V_* = 1 \times 10^{-6}$ m/s, $a = 0.014$, $b = 0.019$, $f_w = 0.3$, and $V_w = 1.0$ m/s. In addition, we continue to assume that normal stress σ is approximately equal to the initial normal stress σ_0 at the location where we interpret the rupture mode. By following the Zheng and Rice analysis summarized above, we can explain the major effect of the prestress level on rupture mode and predict the rupture mode type for different inclination angles.

This application of Zheng and Rice analysis additionally indicates that, for a given angle α , there should be a dependence of the rupture mode on the compressive load P , with higher values of P promoting more crack-like behavior. The analysis predicts that, for the ranges of α and P studied in the presented

theoretical study of velocity-weakening interfaces by Zheng and Rice (21). We also establish experimentally that both pulse-like and crack-like rupture modes can transition to supershear speeds. The resulting supershear rupture speeds are consistent with the analytical predictions of the velocity-weakening model of Samudrala *et al.* (55). The agreement between our experimental observations and models of velocity-weakening faults suggests that velocity-weakening friction plays an important role in dynamic behavior of ruptures and implies that expressing dynamic weakening of friction solely in terms of slip may not be a sufficiently general description.

Some recent numerical studies have pointed out the importance of the rupture initiation process for subsequent rupture dynamics, especially for supershear transition but also for the establishment of

the mode of rupture (9–12, 44). Our current work is focused on modeling the experimental setup numerically to gain further insight into implications of the presented experimental observations, including the role of the nucleation procedure used in experiments. To make a meaningful comparison between simulations and experiments, it is necessary to quantify independently a number of input parameters, including parameters of the rupture nucleation mechanism and friction properties of Homalite interfaces.

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